

Nonlinear orbital stability of traveling wave solutions to an elasto-chemical model*

The equations

Spectral
stability

Exponential
decay of
constructed
 C_0 -semigroup

Nonlinear
orbital stability

Discussion

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2 Spectral stability

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Modelling

- Coupled mechano-chemical system: calcium diffuses freely on a cytoplasmic material (e.g. *cytoskeleton* model of **MURRAY, J. D. AND OSTER, G. F.**, IMA J. Math. Appl. Med. Biol. **1**, pp. 51–75; gel-like substance, elastic properties (deformable)).
- Coupling via: (i) Actomyosin molecules, exert stress on the material; sensitive to calcium concentration; and, (ii) activation of calcium due to deformation of cytoplasm.
- Slow diffusion of calcium: quasi-static balance of forces, inertial terms neglected in the elastic equation.

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- E.g. Post-fertilization traveling waves on eggs. Mechanical and chemical phenomena observed on surface of vertebrate eggs shortly after fertilization. Calcium wave prevents further fertilization.
- Simplest elasto-chemical system, underlies solutions of traveling wave type.

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Model proposed by **D. C. LANE, J. D. MURRAY, AND V. S. MANORANJAN** (IMA J. Math. Appl. Med. Biol. **4** (1987), no. 4, pp. 309–331.)

Chemical model: Calcium released by autocatalytic process (self-estimated), diffuses freely, and more calcium is seen as cytoplasm is stretched (stretch activation)

$$c_t = \underbrace{D\Delta c}_{\text{Fick's diffusion}} + \underbrace{R(c)}_{\text{autocatalytic term}} + \underbrace{\epsilon(\operatorname{div}_x u)}_{\text{stretch activation}}$$

c = free calcium concentration

$\mathbb{R}^3 \ni u$ = elastic displacement vector of cytoplasm

$0 \leq \epsilon$ = stretch factor or contraction stress coefficient

$0 < D$ = Fick's diffusion coefficient of calcium

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Mechanical model: Slow diffusion of calcium, quasi-static balance of forces. Inertial terms negligible!

$$\nabla \cdot \sigma + \rho F = 0,$$

$$\begin{aligned} \sigma = & \underbrace{\mu_1 e_t + \mu_2 (\operatorname{div}_x u)_t I}_{\text{viscous stress}} + \underbrace{E e + (1 - E)(\operatorname{div}_x u) I}_{\text{elastic stress}} + \\ & + \underbrace{\tau(c) I}_{\text{active contraction stress due to calcium}}, \end{aligned}$$

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Body forces proportional to elastic displacement:

$$F = -su, \quad s > 0,$$

$$e = \frac{1}{2}(\nabla u + \nabla u^T) = \text{strain tensor}$$

$\rho =$ cytogel density
 $\mu_i =$ bulk and shear viscosities
 $E =$ elastic modulus

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System of equations

One dimensional version of the model by **LANE** *et al.* :

$$\mu u_{xxt} + u_{xx} - \tau(c)_x - su = 0,$$

$$c_t - Dc_{xx} - R(c) - \epsilon u_x = 0,$$

$(x, t) \in \mathbb{R} \times [0, +\infty)$, where:

$\mathbb{R} \ni u =$ elastic displacement,

$\mathbb{R} \ni c =$ concentration of free calcium,

$\mu = \mu_1 + \mu_2 =$ combined shear and bulk viscosities,

$0 < s =$ restoring force,

$0 \leq \epsilon =$ contraction stress on the increase of c ,

$0 < D =$ Fick's diffusion constant of calcium.

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Nonlinear terms:

$R(c)$ = autocatalytic term,

$\tau(c)$ = contractile forces acting on the medium due to c .

Assumptions:

1. $R(0) = R(1) = R(c_0) = 0$, for some $c_0 \in (0, 1)$.
2. $\int_0^1 R(c) dc > 0$.
3. $R'(0) < 0$, and $R'(1) < 0$ (bistable shape),
4. $\tau'(c)$ bounded for all $c \in (0, 1)$,
5. τ has compact support in $(0, 1)$, with $\tau \equiv 1$ for $c \in (\delta, 1 - \delta)$, $1 \gg \delta > 0$.

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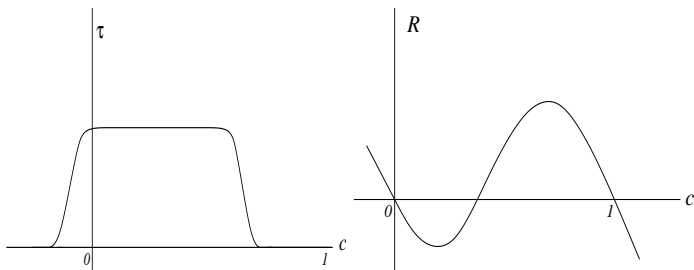
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Form of the nonlinear terms $R(c)$ and $\tau(c)$.

Traveling wave solutions

$$(u, c)(x, t) = (\bar{u}, \bar{c})(x + \theta t),$$

$\theta =$ wave speed,

with

$$\bar{u}(\pm\infty) = 0, \quad \bar{c}(+\infty) = 1, \quad \bar{c}(-\infty) = 0.$$

Wave equations:

$$\mu\theta\bar{u}'''' + \bar{u}'' - (\tau(\bar{c}))' - s\bar{u} = 0,$$

$$\theta\bar{c}' - D\bar{c}'' - R(\bar{c}) - \epsilon\bar{u}' = 0.$$

$x \rightarrow x + \theta t, ' = d/dx$ (galilean variable).

Existence of traveling waves for small $\epsilon \geq 0$

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- **G. FLORES, A. MINZONI, K. MISCHIAKOW, AND V. MOLL**, *Nonlinear Anal.* **36** (1999), no. 1, Ser. A: Theory Methods, pp. 45–62.

Proposition

For $\epsilon \geq 0$ suff. small, there exist $(\bar{u}^\epsilon, \bar{c}^\epsilon)$ such that $\bar{u}^\epsilon(\pm\infty) = 0$, $\bar{c}^\epsilon(+\infty) = 1$, $\bar{c}^\epsilon(-\infty) = 0$, and the speed is uniquely determined by

$$\theta(\epsilon) = \theta_0 + o(1),$$
$$\theta_0 := \frac{\int_0^1 R(c) dc}{\int_{\mathbb{R}} \bar{c}'(x)^2 dx} > 0,$$

When $\epsilon = 0$, \bar{c}^0 is the bistable Nagumo front, with speed θ_0 .

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Corollary (Exponential decay)

$$\begin{aligned} |\partial_x^j \bar{u}^\epsilon(x)| &\lesssim e^{-|x|/C_1}, \quad \text{as } |x| \rightarrow +\infty, \quad j = 0, 1, 2, \\ |\partial_x^i (\bar{c}^\epsilon(x) - 1)| &\lesssim e^{-x/C_1}, \quad \text{as } x \rightarrow +\infty, \quad i = 0, 1, \\ |\partial_x^i \bar{c}^\epsilon(x)| &\lesssim e^{+x/C_1}, \quad \text{as } x \rightarrow -\infty, \quad i = 0, 1, \end{aligned}$$

uniform $C_1 > 0$, for all $\epsilon \sim 0^+$.

Perturbed problem

Let $u + \bar{u}$, $c + \bar{c}$ be solutions, with (u, c) perturbations.

Nonlinear perturbation equations

$$\begin{aligned}\mu\theta u_{xxx} + \mu u_{xxt} + u_{xx} - su - (\tau(c + \bar{c}) - \tau(\bar{c}))_x &= 0, \\ c_t + \theta c_x - Dc_{xx} - \epsilon u_x - (R(c + \bar{c}) - R(\bar{c})) &= 0.\end{aligned}\tag{NL}$$

Linearized (around the waves) system for the perturbation

$$\begin{aligned}\mu\theta u_{xxx} + \mu u_{xxt} + u_{xx} - su - (\tau'(\bar{c})c)_x &= 0, \\ c_t + \theta c_x - Dc_{xx} - \epsilon u_x - R'(\bar{c})c &= 0.\end{aligned}\tag{L}$$

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Results

- **G. FLORES AND R. P.**, Journal of Differential Eqs. **247** (2009), no. 5, pp. 1529–1590.

Theorem 1 (Spectral stability)

For each $\epsilon \geq 0$ sufficiently small, traveling waves $(\bar{u}^\epsilon, \bar{v}^\epsilon)$ are spectrally stable and $\lambda = 0$ is an isolated simple eigenvalue.

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Theorem 1 (Spectral stability)

For each $\epsilon \geq 0$ sufficiently small, traveling waves $(\bar{u}^\epsilon, \bar{u}^\epsilon)$ are spectrally stable and $\lambda = 0$ is an isolated simple eigenvalue.

Theorem 2 (Semigroup estimates)

For each $\epsilon \geq 0$ sufficiently small, there exists $\omega_0 > 0$ such that for each inicial cond. $(u_0, c_0) \in H^2 \times H^1$ there is a global solution (u_x, c) en $C([0, +\infty); H^1 \times H^1)$ to system (L) and some $\alpha_ \in \mathbb{R}$ such that*

$$\|(u_x, c)(\cdot, t) - \alpha_* (\bar{u}_{xx}^\epsilon, \bar{c}_x^\epsilon)(\cdot)\|_{L^2 \times L^2} \lesssim e^{-\omega_0 t},$$

for each $t > 0$. Moreover, if $(u_0, c_0) \in (W^{1,1} \cap H^3) \times H^2$ then

$$\|(u, c)(\cdot, t) - \alpha_* (\bar{u}_x^\epsilon, \bar{c}_x^\epsilon)(\cdot)\|_{L^\infty \times L^\infty} \rightarrow 0, \quad \text{if } t \rightarrow +\infty.$$

Theorem 3 (Nonlinear stability)

For each $\epsilon \geq 0$ sufficiently small there exists $\eta_0 > 0$ such that, if $(\tilde{u}_0, \tilde{c}_0) \in H^2 \times H^1$ and $\alpha_0 \in \mathbb{R}$ satisfy

$$\|(\tilde{u}_{0x}, \tilde{c}_0)(\cdot) - (\bar{u}_x^\epsilon, \bar{c}^\epsilon)(\cdot + \alpha_0)\|_{H^1 \times H^1} < \eta \leq \eta_0,$$

Then there exists a unique global solution

$(\tilde{u}_x, \tilde{c}) \in C([0, +\infty); H^1 \times H^1)$ to system (NL) and some $\alpha_\infty \in \mathbb{R}$ such that

$$|\alpha_0 - \alpha_\infty| < C_1 \eta_0,$$

$$\|(\tilde{u}_x, \tilde{c})(\cdot, t) - (\bar{u}_x^\epsilon, \bar{c}^\epsilon)(\cdot + \theta t + \alpha_\infty)\|_{H^1 \times H^1} \leq C \eta_0 e^{-\frac{1}{2}\omega_0 t} \rightarrow 0,$$

as $t \rightarrow +\infty$. Moreover, if $(\tilde{u}_0, \tilde{c}_0) \in (W^{1,1} \cap H^3) \times H^2$ then

$$\|(u, c)(\cdot, t) - (\bar{u}^\epsilon, \bar{c}^\epsilon)(\cdot + \theta t + \alpha_\infty)\|_{L^\infty \times L^\infty} \leq C \eta_0 e^{-\omega_0 t} \rightarrow 0,$$

as $t \rightarrow +\infty$.

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Spectral problem

Perturbations of form $(e^{\lambda t}u(x), e^{\lambda t}c(x))$ with $\lambda \in \mathbb{C}$:

$$\begin{aligned}\mu\theta u_{xxx} + (\mu\lambda + 1)u_{xx} - su - (\tau'(\bar{c})c)_x &= 0, \\ \lambda c + \theta c_x - Dc_{xx} - \epsilon u_x - R'(\bar{c})c &= 0.\end{aligned}$$

Necessary condition for stability: no solutions $(u, c) \in L^2$ with $\operatorname{Re} \lambda > 0$.

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First order system formulation

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Reine Angew. Math. **410** (1990), pp. 167–212.

Spectral problem:

$$W_x = \mathbb{A}^\epsilon(x, \lambda)W,$$

$$W := (u, u_x, u_{xx}, c, c_x)^\top.$$

$$W \in L^2(\mathbb{R}; \mathbb{C}^5)$$

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$$\mathbb{A}^\epsilon(x, \lambda) := \begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ s/\mu\theta & 0 & -(1 + \mu\lambda)/\mu\theta & \tau''(\bar{c})\bar{c}_x/\mu\theta & \tau'(\bar{c})/\mu\theta \\ 0 & 0 & 0 & 0 & 1 \\ 0 & -\epsilon/D & 0 & (\lambda - R'(\bar{c}))/D & \theta/D \end{pmatrix}$$

Asymptotic systems:

$$W_x = \mathbb{A}_\pm^\epsilon(\lambda)W,$$

$$\mathbb{A}^\epsilon(\lambda) :=$$

$$= \begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ s/\mu\theta & 0 & -(1 + \mu\lambda)/\mu\theta & 0 & \tau'(n)/\mu\theta \\ 0 & 0 & 0 & 0 & 1 \\ 0 & -\epsilon/D & 0 & (\lambda - R'(n))/D & \theta/D \end{pmatrix}$$

$$n = 0, 1 \text{ for } x = -\infty, +\infty.$$

Definition of spectra

Family of densely defined closed operators in $L^2(\mathbb{R}; \mathbb{C}^5)$:

$$\begin{aligned} \mathcal{T}(\lambda) &: \mathcal{D}(\mathcal{T}) \longrightarrow L^2(\mathbb{R}; \mathbb{C}^5), \\ \mathcal{T}(\lambda)W &:= W_x - \mathbb{A}^\epsilon(x, \lambda)W, \end{aligned}$$

domain $\mathcal{D}(\mathcal{T}^\epsilon) = H^1(\mathbb{R}; \mathbb{C}^5)$, indexed by $\epsilon \geq 0$ and $\lambda \in \mathbb{C}$.

Definition (Spectra)

$$\rho := \{\lambda \in \mathbb{C} : \mathcal{T}(\lambda) \text{ is 1 - 1 and onto } \mathcal{T}(\lambda)^{-1} \text{ bounded}\},$$

$$\sigma_{pt} := \{\lambda \in \mathbb{C} : \mathcal{T}(\lambda) \text{ is Fredholm with index 0 and} \\ \text{non-trivial kernel}\},$$

$$\sigma_{ess} := \{\lambda \in \mathbb{C} : \mathcal{T}(\lambda) \text{ either has non-zero index or is not Fredholm}\}$$

The spectrum is $\sigma = \sigma_{ess} \cup \sigma_{pt}$. Since \mathcal{T} closed, then $\rho = \mathbb{C} \setminus \sigma$. We say $\lambda \in \sigma_{pt}$ is an eigenvalue.

Definition (Multiplicities)

For $\lambda \in \sigma_{pt}$: (i) Its geometric multiplicity (g.m.) is the maximal number of linearly independent elements in $\ker \mathcal{T}(\lambda)$. (ii) Suppose $\lambda \in \sigma_{pt}$ has g.m. = 1, so that $\ker \mathcal{T}(\lambda) = \text{span} \{W_1\}$. We say λ has algebraic multiplicity (a.m.) equal to m if we can solve

$$\mathcal{T}^\epsilon(\lambda)W_j = \tilde{\mathbb{A}}_1^\epsilon(x)W_{j-1},$$

for each $j = 2, \dots, m$, with $W_j \in H^1$, but there is no H^1 solution W to

$$\mathcal{T}^\epsilon(\lambda)W = \tilde{\mathbb{A}}_1^\epsilon(x)W_m.$$

For an arbitrary eigenvalue $\lambda \in \sigma_{pt}$ with g.m. = l , the a.m. is defined as the sum of the multiplicities $\sum_k^l m_k$ of a maximal set of linearly independent elements in $\ker \mathcal{T}(\lambda) = \text{span} \{W_1, \dots, W_l\}$.

Remark : The definition coincides with the usual one for equations in standard form $U_t = LU$, for a given linearized operator L , when written as a first order system. This holds because *the Fredholm properties of $\mathcal{L} - \lambda$ and $\mathcal{T}(\lambda)$ are the same.*
See:

- B. SANDSTEDE, *Stability of travelling waves*, in Handbook of dynamical systems, Vol. 2, B. Fiedler, ed., North-Holland, Amsterdam, 2002, pp. 983–1055.
- B. SANDSTEDE AND A. SCHEEL, Proc. Roy. Soc. Edinburgh Sect. A **130** (2000), no. 2, pp. 419–448.
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Definition (Spectral stability)

We say the waves are spectrally stable if

$$\sigma \subset \{\lambda \in \mathbb{C} : \operatorname{Re} \lambda < 0\} \cup \{0\},$$

i.e., there are no solutions in L^2 with $\operatorname{Re} \lambda \geq 0$; here $\lambda = 0$ is the eigenvalue associated to translation invariance, with eigenfunction (\bar{u}_x, \bar{c}_x) .

Idea:

- For $\epsilon = 0$, spectral stability of the Nagumo and elastic fronts follow by energy estimates.
- For $\epsilon \geq 0$ sufficiently small, stability persists due to uniform convergence of the *Evans functions* as $\epsilon \rightarrow 0^+$.

Evans function

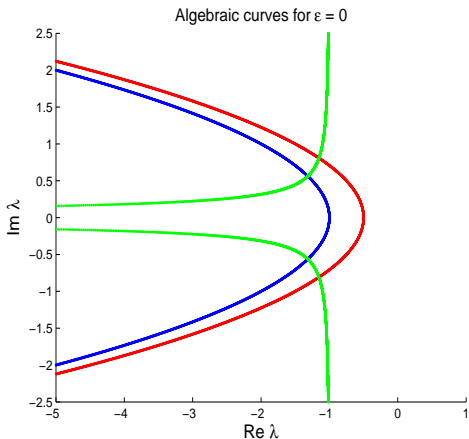
- $\Omega =$ open region in the complement of the essential spectrum containing $\{\operatorname{Re} \lambda \geq 0\}$.
- The Evans function $D^\epsilon(\lambda)$ is an analytical function defined on Ω ; its zeroes coincide in location and multiplicity with the eigenvalues of the spectral problem.
- D can be defined via the Wronskian of the first order system.

$$\Lambda_0 := \min\{|R'(1)|, |R'(0)|, 1/\mu\} > 0$$

$$\Omega := \{\lambda \in \mathbb{C} : \operatorname{Re} \lambda > -\frac{1}{2}\Lambda_0\}.$$

In Ω the dimensions of the stable and unstable manifolds of the asymptotic systems $W_x = \mathbb{A}_\pm W$ are constant:

$$\dim U_\pm^\epsilon(\lambda) = 2, \quad \dim S_\pm^\epsilon(\lambda) = 3.$$



Algebraic curves limiting the essential spectrum for $\epsilon \sim 0^+$. Note the spectral gap.

Definition of Evans function:

$$D^\epsilon(\lambda) = \det \left(W_1^-(x, \lambda), W_2^-(x, \lambda), W_3^+(x, \lambda), \right. \\ \left. W_4^+(x, \lambda), W_5^+(x, \lambda) \right) |_{x=0},$$

Properties: D^ϵ analytic in Ω , and $D^\epsilon = 0$ iff λ is an eigenvalue.
The order of the zero coincides with the algebraic multiplicity.

Facts:

- (\bar{u}^0, \bar{c}^0) are spectrally stable (proof with energy estimates). $\lambda = 0$ is a simple eigenvalue, i.e., $D^0(\lambda) \neq 0$ for $\operatorname{Re} \lambda \geq 0$, except for $\lambda = 0$; $(d/d\lambda)D^0(0) \neq 0$.
- For all $\epsilon \sim 0^+$, $\lambda \in \Omega$, $\mathbb{A}_{\pm}^{\epsilon}(\lambda)$ are hyperbolic with $\dim U_{\pm}^{\epsilon}(\lambda) = 2$, $\dim S_{\pm}^{\epsilon}(\lambda) = 3$. This shows $\sigma_{\text{ess}} \subset \{\operatorname{Re} \lambda < 0\}$.
- For all $\epsilon \geq 0$ fixed, $\sigma_{\text{pt}} \cap \{\operatorname{Re} \lambda \geq -1/2\mu\}$ is uniformly bounded.

- For $\lambda \in \Omega$, $\epsilon \geq 0$ small, $U_-^\epsilon, S_+^\epsilon \rightarrow U_-^0, S_+^0$ in angle as $\epsilon \rightarrow 0^+$ with rate $\mathcal{O}(\epsilon + \delta(\epsilon)) =: \eta(\epsilon)$, i.e.,

$$|v_{j\pm}^\epsilon - v_{j\pm}^0| \leq \eta(\epsilon).$$

for spanning bases. Moreover, by exp. decay of the waves,

$$|(\mathbb{A}^\epsilon - \mathbb{A}_\pm^\epsilon) - (\mathbb{A}^0 - \mathbb{A}_\pm^0)| \leq C_2 \eta(\epsilon) e^{-|x|/C_1}.$$

Convergence result:

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Under such structural conditions (hiperbolicity, exponential decay, regularity) the sequence of Evans functions D^ϵ converges locally and uniformly to D^0 as $\epsilon \rightarrow 0^+$ with ratio

$$|D^\epsilon - D^0| \leq \mathcal{O}(\epsilon + |\delta(\epsilon)|) = \mathcal{O}(\eta(\epsilon));$$

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$D^0(\lambda) \neq 0$ in $\operatorname{Re} \lambda \geq 0$, except at $\lambda = 0$. Thus, $D^\epsilon(\lambda) \neq 0$ for $\epsilon \sim 0$ small (by uniform convergence and analyticity). There are no non-zero eigenvalues with $\{\operatorname{Re} \lambda \geq 0\}$. The multiplicity persists by convergence of the derivatives of D^ϵ .

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Key ideas:

1. Write the linear system (L) in terms of the deformation gradient

$$(v, c) := (u_x, c),$$

$$\mathcal{L}(\partial_x, \partial_t)(v, c) = 0.$$

2. Construction of the semigroup associated to the above equation. Global solution operator:

$$(v_0, c_0) \in H^1 \times H^1 \mapsto \mathcal{S}(t)(v_0, c_0) \in C([0, +\infty); H^1 \times H^1).$$

Lemma

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Lemma

$\mathcal{S}(t)$ is a C_0 -semigroup in $H^1 \times H^1$.

3. $\mathcal{S}(t)$ has a densely defined generator \mathcal{A}

$$\mathcal{A} : \mathcal{D} \subset H^1 \times H^1 \longrightarrow H^1 \times H^1,$$

where

$$\mathcal{D} = (H^2 \cap L^1 \cap \mathcal{U}) \times H^3,$$

and

$$\mathcal{U} = \{u \in L^2 : \int u = 0, \int \int^x u = 0, \text{ two antiderivatives in } L^2\}.$$

($\hat{u}(k)$ has a double zero in $k = 0$.)

4. Fredholm properties:

$$\begin{aligned}\text{nul}(\mathcal{A} - \lambda) &\leq \text{nul}(\mathcal{T}(\lambda)), \\ \text{nul}(\mathcal{A}^* - \lambda^*) &\leq \text{nul}(\mathcal{T}(\lambda)^*),\end{aligned}$$

imply

$$\begin{aligned}\sigma_{\text{pt}}(\mathcal{A}) &\subseteq \sigma, \\ \sigma_{\text{ess}}(\mathcal{A}) &\subseteq \sigma,\end{aligned}$$

Thus, spectral stability for small $\epsilon \geq 0$.

5. $(\bar{u}_{xx}, \bar{c}_x)$ – eigenfunction associated to $\lambda = 0$. $(\bar{\psi}, \bar{\phi})$ – adjoint eigenfunction.

Projection:

$$\mathcal{P}_1(v, c) := (v, c) - \frac{\langle (v, c), (\bar{\psi}, \bar{\phi}) \rangle_{H^1 \times H^1}}{\Theta} (\bar{u}_{xx}, \bar{c}_x),$$

$$\Theta := \langle (\bar{u}_{xx}, \bar{c}_x), (\bar{\psi}, \bar{\phi}) \rangle_{H^1 \times H^1} \neq 0.$$

$$X_1 := \text{rank of } \mathcal{P}_1 \subset H^1 \times H^1,$$

$$\mathcal{A}_1 := \mathcal{A}|_{X_1}$$

$$\mathcal{S}_1(t) := \mathcal{S}(t)\mathcal{P}_1$$

is a C_0 -semigroup in X_1 .

5. Resolvent estimates:

(a) For each $\epsilon \geq 0$, $\operatorname{Re} \lambda \geq 0$ and $|\operatorname{Im} \lambda|$ suff. big,

$$\|(\lambda - \mathcal{A})^{-1}\|_{H^1 \rightarrow H^1} \leq C,$$

(b) For $\epsilon \geq 0$ suff. small and $\operatorname{Re} \lambda \geq 0$ big enough

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6. Gearhart-Prüss criterion (C_0 -semigroups in Hilbert spaces):

Since

$$\sup\{\operatorname{Re} \lambda : \lambda \in \sigma(\mathcal{A}_1)\} < 0,$$

$$\sup_{\operatorname{Re} \lambda > 0} \|(\lambda - \mathcal{A}_1)^{-1}\|_{X_1 \rightarrow X_1} < +\infty,$$

then the semigroup $\mathcal{S}_1(t)$ on the Hilbert space X_1 is *exponentially stable*:

$$\|\mathcal{S}_1(t)(v, c)\|_{H^1} \leq Ce^{-\omega_0 t} \|(v, c)\|_{H^1}, \quad (v, c) \in X_1.$$

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Ideas:

1. Ansatz of form:

$$(\tilde{v}, \tilde{c})(x, t) = (v, c)(x + \theta t + \alpha(t), t) + (\bar{u}_x, \bar{c})(x + \theta t + \alpha(t)),$$

$\alpha(t)$ = modulated phase depending on t .

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Implicit function theorem applied to the functional

$$\mathcal{G}[(\tilde{v}, \tilde{c}), \alpha](t) := \langle (v, c), (\bar{\psi}, \bar{\phi}) \rangle_{H^1} = 0$$

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$$\dot{\alpha}(t) = \frac{\Theta^{-1} \langle (\mu^{-1} N_1, N_2), (\bar{\psi}, \bar{\phi}) \rangle_{H^1 \times H^1}}{1 - \Theta^{-1} \langle (v, c), (\bar{\psi}_\xi, \bar{\phi}_\xi) \rangle_{H^1 \times H^1}}$$

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$$\|(\tilde{v}, \tilde{c})(\cdot, t) - (\bar{u}_x, \bar{c})(\cdot + \theta t + \alpha(t))\|_{H^1 \times H^1} \leq C\eta_0 e^{-\frac{1}{2}\omega_0 t} \rightarrow 0,$$

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Details in: FLORES-P, Journal of Differential Eqs. **247** (2009), no. 5, pp. 1529–1590.

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Thanks!