

Mapeo estándar

Sabemos que el flujo de las ecs. de Hamilton es un simplectomorfismo $\phi_{t, t_0}^L: \mathbb{R}^{2n} \rightarrow \mathbb{R}^{2n}$

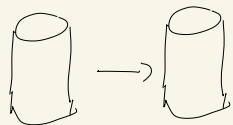
En general, no es posible escribir una fórmula explícita para ϕ .

$$f: \mathbb{R}^2 \rightarrow \mathbb{R}^2$$

$$\begin{bmatrix} q \\ p \end{bmatrix} \mapsto \begin{bmatrix} q + p + \varepsilon G(q) \\ p + \varepsilon G'(q) \end{bmatrix}$$

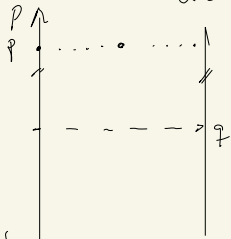
$$G(q) = \frac{1}{2\pi} \sin(2\pi q) \quad \text{función periódica de periodo 1}$$

$$f: \mathcal{S}^1 \times \mathbb{R} \rightarrow \mathcal{S}^1 \times \mathbb{R}$$



si $p \in \mathbb{Q}$ $p = \frac{l}{m}$

$l \in \mathbb{Z}, m \in \mathbb{N}$
 la órbita es periódica.
 si $p \in \mathbb{R}/\mathbb{Q}$ la órbita llena densamente el círculo.



Caso $\varepsilon = 0$

$$f \begin{pmatrix} q \\ p \end{pmatrix} = \begin{pmatrix} q + p \\ p \end{pmatrix}$$

Todos los círculos $p = \text{cte}$ son invariantes bajo la dinámica.

$\varepsilon \neq 0$ dinámica más complicada.

$$\begin{pmatrix} q + p + \frac{\varepsilon}{2\pi} \sin(2\pi q) \\ p + \frac{\varepsilon}{2\pi} \sin(2\pi q) \end{pmatrix} = \begin{pmatrix} q + p + \varepsilon G(q) \\ p + \varepsilon G'(q) \end{pmatrix}$$

Veamos que el mapeo es simpléctico.

$$Df \begin{pmatrix} q \\ p \end{pmatrix} = \begin{pmatrix} 1 + \varepsilon G'(q) & 1 \\ \varepsilon G'(q) & 1 \end{pmatrix}$$

$$Df \begin{pmatrix} q \\ p \end{pmatrix}^T J Df \begin{pmatrix} q \\ p \end{pmatrix} =$$

$$= \begin{pmatrix} 1 + \varepsilon G'(q) & \varepsilon G'(q) \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 1 + \varepsilon G'(q) & 1 \\ \varepsilon G'(q) & 1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 + \varepsilon G'(q) & \varepsilon G'(q) \\ 1 & 1 \end{pmatrix} \begin{pmatrix} \varepsilon G'(q) & 1 \\ -(1 + \varepsilon G'(q)) & -1 \end{pmatrix}$$

$$= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = J$$

El mapeo estándar es simpléctico.

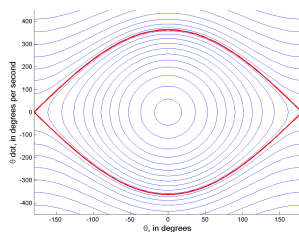
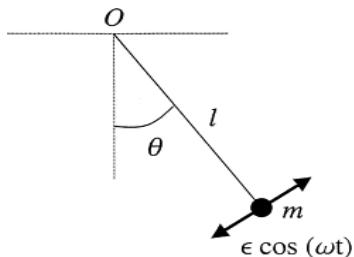
$$\det(Df \begin{pmatrix} q \\ p \end{pmatrix}) = 1 + \varepsilon G'(q) - \varepsilon G'(q) = 1$$



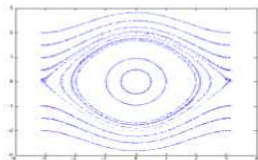
Péndulo perturbado

$$\dot{x} = y$$

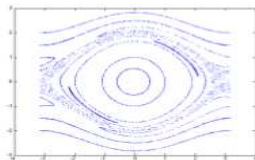
$$\dot{y} = -\omega^2 \sin(x) + \epsilon \cos(\omega t)$$



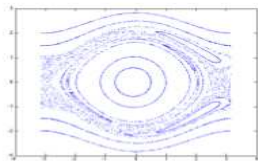
El espacio fase se “destruye”.



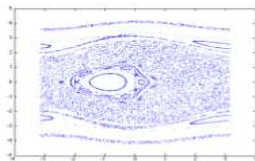
(a) $\epsilon = 0$



(b) $\epsilon = 0,01$



(c) $\epsilon = 0,1$



(d) $\epsilon = 1,0$

Objetos invariantes en mapeos

- ▶ Puntos fijos

$$\phi_{t,t_0}(x) \quad x \in (\mathbb{R}, \mathbb{P})$$

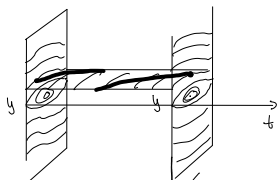
$$f(x) = x$$

- ▶ Órbitas periódicas

$$f^n(x) = x$$

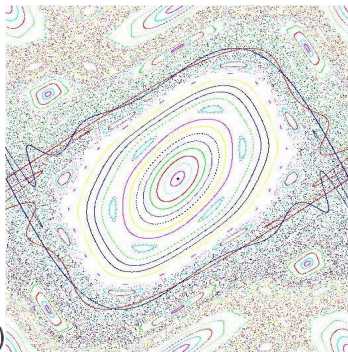
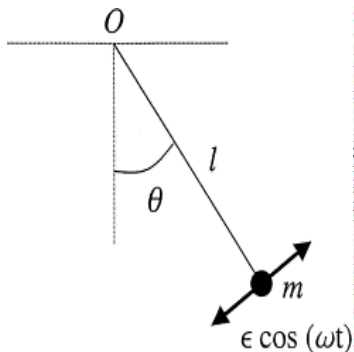
- ▶ Círculos invariantes

$$f(\text{círculo}) = \text{círculo}$$



Péndulo perturbado

$$H(y, x) = \frac{1}{2}y^2 + \varepsilon V(x) \times \sum_{n=-\infty}^{\infty} \delta\left(\frac{t}{T} - n\right)$$

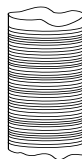


Puntos fijos, órbitas periódicas y círculos invariantes

Mapeo integrable.

$$y' = y$$

$$x' = x + \omega(y) \pmod{1}$$

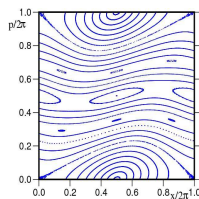


ω es el número de rotación.

Mapeo estándar.

$$y' = y + \frac{\varepsilon}{2\pi} \sin(2\pi x)$$

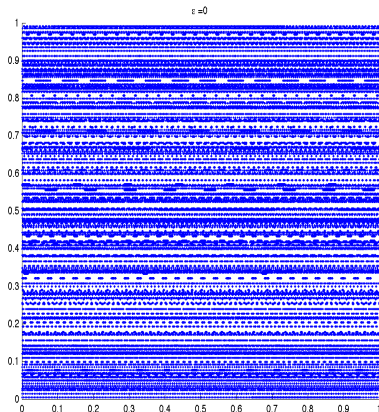
$$x' = x + y' \pmod{1}$$



Moviendo ε

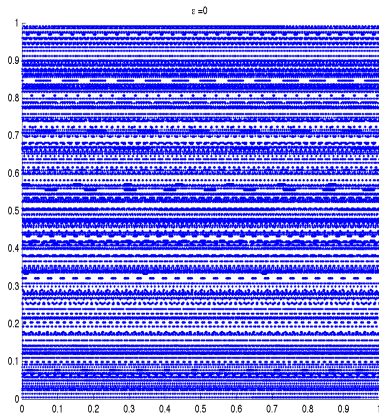
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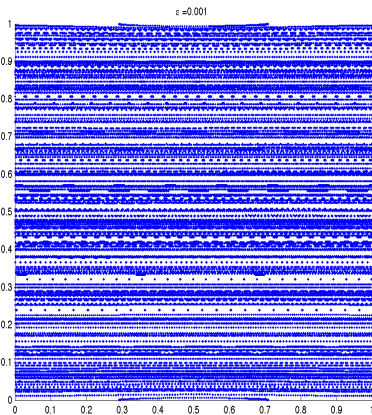
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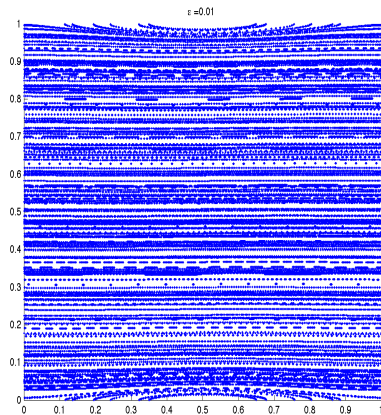
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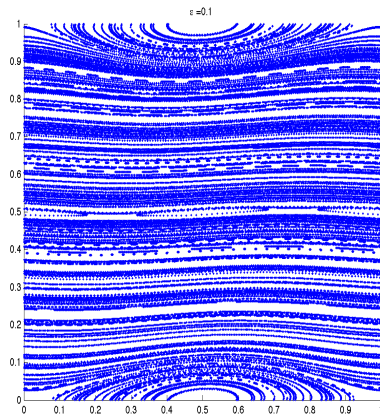


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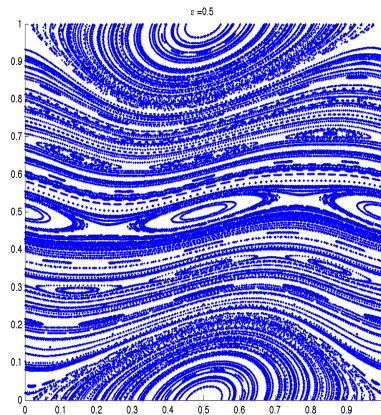
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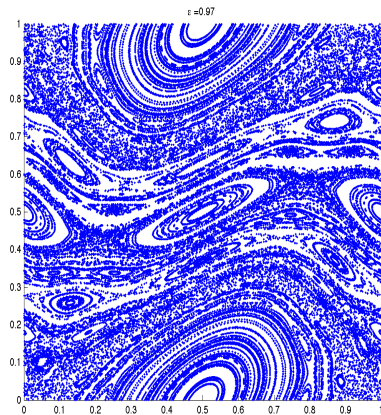
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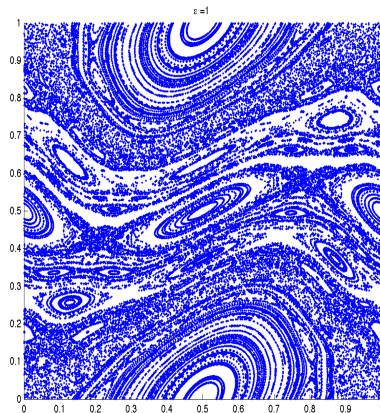
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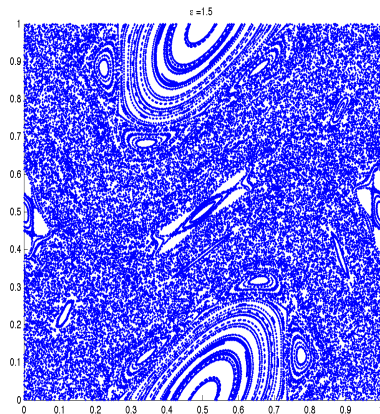
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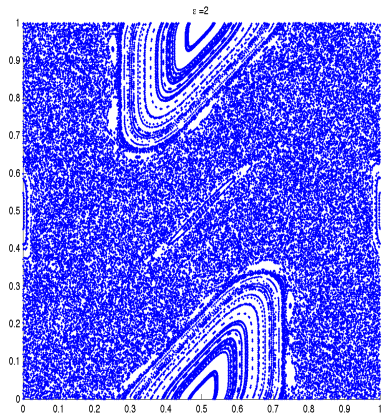
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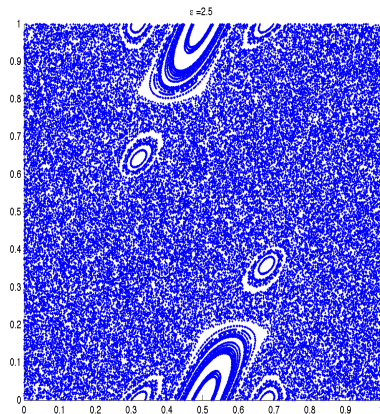
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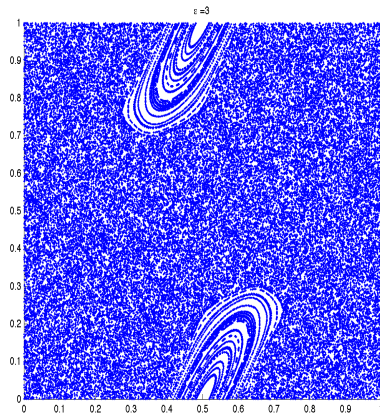
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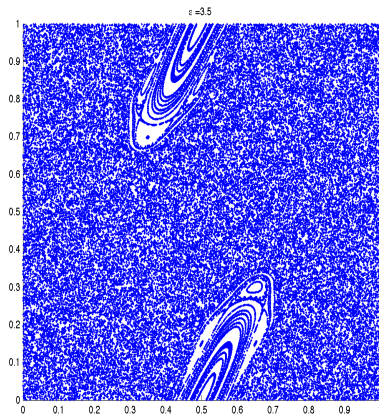
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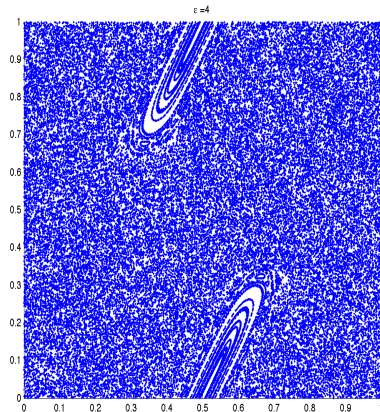
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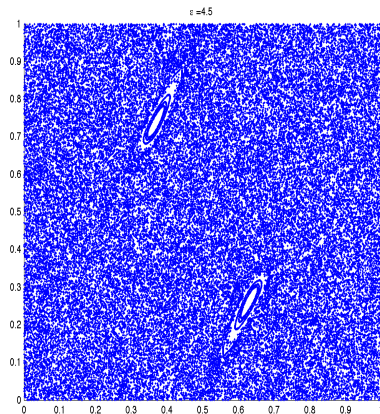
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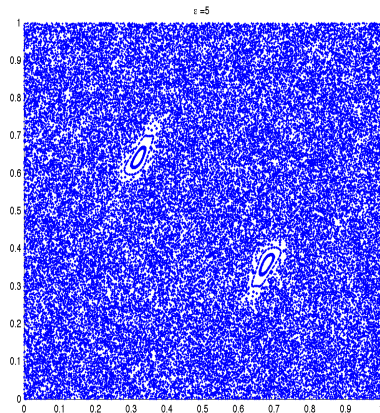
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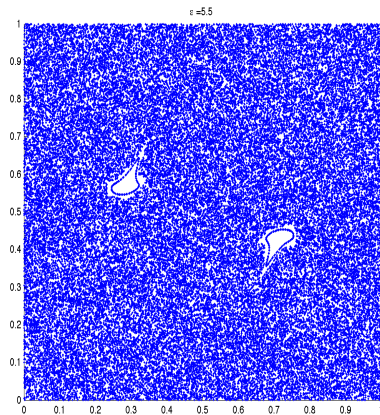
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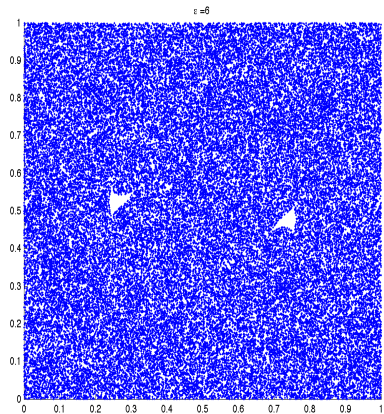
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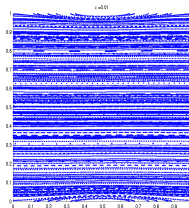
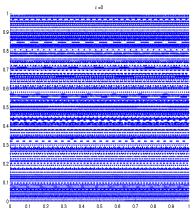


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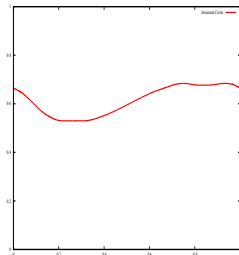
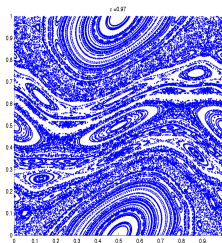
Teorema KAM

Las órbitas quasi-periódicas persisten cerca del caso integrable



Teorema [C. - de la Llave]

Las órbitas quasi-periódicas persisten lejos del caso integrable. Desaparecen cuando la parametrización pierde regularidad (de Sobolev).



Además la prueba del teorema sugiere algoritmos muy eficientes ($\mathcal{O}(N \log N)$ operaciones).