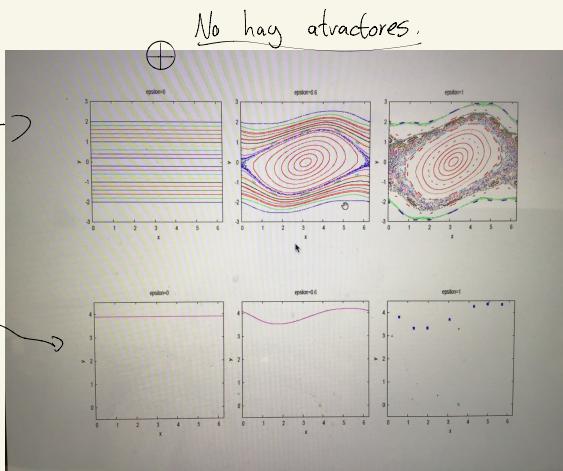
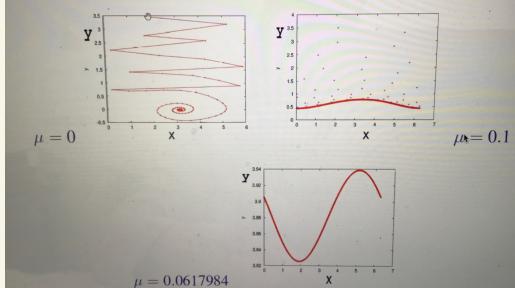


Caso conservativo

Caso dissipativo



Looking for the torus with frequency $\omega = 2\pi \frac{\sqrt{5}-1}{2} \simeq 3.8832$, dissipative standard map with $\varepsilon = 0.1$, $\lambda = 0.9$.



Este es un buen ejercicio!

En el caso dissipativo sólo se tiene un toro invariante.

Modeling the spin-orbit problem

- Spin-orbit problem:
 - ▷ triaxial satellite \mathcal{S} (with $I_1 < I_2 < I_3$);
 - ▷ center of the satellite moves on a **Keplerian orbit** around a central planet \mathcal{P} ;
 - ▷ spin-axis perpendicular to orbit plane and coinciding with **shortest physical axis**;
 - ▷ NO tidal torque, due to the non-rigidity.

- Conservative case: equation of motion:

$$\ddot{x} + \varepsilon \left(\frac{a}{r(t)} \right)^3 \sin(2x - 2f(t)) = 0, \quad \varepsilon = \frac{3}{2} \frac{I_2 - I_1}{I_3}$$

corresponding to a 1-dim, time-dependent Hamiltonian:

$$\mathcal{H}(y, x, t) = \frac{y^2}{2} - \frac{\varepsilon}{2} \left(\frac{a}{r(t)} \right)^3 \cos(2x - 2f(t)).$$

Spin-orbit problem

- $\gamma_g^\pm = \frac{\sqrt{5} \pm 1}{2}$

- Diophantine frequency: $\omega = \frac{1}{2 + \gamma_g^\pm}$ or $\omega = \gamma_g^\pm$.

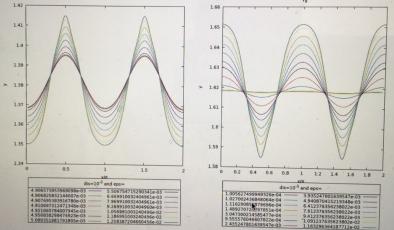
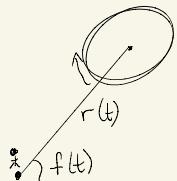


Figure: The invariant attractor for SO with $\eta = 10^{-3}$. Left: ω . Right: γ_g^+ .

Modelo dissipativo



P. Goldreich, "Final spin states of planets and satellites," *Astronom.J.*, 71, no 6, 425-438 (1966)
 S. J. Peale, "Icarus 78 4-18 (2005)"

Torca de marea

$$\mathcal{T}(z, t) = -K_d [L(e, t)z - N(e, t)]$$

$$L(e, t) = \frac{\alpha^6}{r(t)^6}, \quad N(e, t) = \frac{\alpha^6}{r(t)^6} f(t)$$

K_d - constante dissipativa
número de Lava (estructura interior del cuerpo)

$$K_d = 3\pi \frac{k_2}{Q} \left(\frac{R_e}{\alpha}\right)^3 \frac{M}{M_{\odot}}$$

Re - radio ecuatorial
Mea de P
Mea de S
factor de calidad

movimiento medio constante de velocidad, $\bar{J}_2 = \frac{3}{2} \pi R_e^2$
 $K_d \approx 10^{-8}$. En astronomía las fuerzas dissipativas son
pequeñas pero actúan por tiempos muy
longos.

A.C.M Correia, J. Laskar, "Mercury's capture into 3/2 spin-orbit resonance as a result of its chaotic dynamics," *Nature* 429, 848-850 (2004)

Tomar promedios de $L(e, t)$, $N(e, t)$

$$\bar{\mathcal{T}} = \overline{\mathcal{T}}(z) = -K_d [\bar{L}(e)z - \bar{N}(e)]$$

Tarea Repetir la derivación del mapeo
estándar pero con las torcas promediadas

$$\dot{r} = \bar{J}$$

$$\dot{\bar{J}} = -\varepsilon \cos(\bar{J}, e) \sin 2\bar{J} - K_d [\bar{L}(e)\bar{z} - \bar{N}(e)]$$

Sin dissipación \rightarrow mapeo estándar

Con dissipación \rightarrow deducir el mapeo correspondiente.